

The all-electron full-potential linearized augmented plane-wave method

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Outline

Theoretical background

- The FLAPW method and the LAPW basis
- Separation of core electrons from valence electrons
- Representation of density and potential

FLAPW in practice

- Setting the parameters
- Semicore states and ghost bands
- The linearization error

Using fleur

- The input file generator
- The inp.xml file
- Further reading
- Conclusion





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Motivation: FLAPW in zoo of electronic structure methods













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MAX DRIVING THE EXASCALE TRANSITION

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The LAPW basis

 Atom-centered functions in MT spheres matched in value and slope to plane waves in interstitial region (IR)

$$\phi_{\mathbf{kG}}(\mathbf{r}) = \begin{cases} \frac{1}{\sqrt{\Omega}} e^{j(\mathbf{k}+\mathbf{G})\mathbf{r}} & \text{for } \mathbf{r} \in \mathbf{IR} \\ \sum_{L} \left[a_{\mathbf{kG}}^{L\alpha} u_{l}^{\alpha}(r_{\alpha}, E_{l}^{\alpha}) + b_{\mathbf{kG}}^{L\alpha} \dot{u}_{l}^{\alpha}(r_{\alpha}, E_{l}^{\alpha}) \right] Y_{L}(\hat{\mathbf{r}}_{\alpha}) & \text{for } \mathbf{r} \in \mathbf{MT}^{\alpha} \end{cases}$$

• u_l^{α} and \dot{u}_l^{α} are solutions and energy derivatives for the spherical potential at energy parameters E_l^{α}







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• u_l^{α} and \dot{u}_l^{α} are solutions and energy derivatives for the spherical potential at energy parameters E_l^{α}

Parameters:

• $K_{\max} = |\mathbf{k} + \mathbf{G}|_{\max}$ reciprocal plane wave cutoff • I_{\max}^{α} angular momentum cutoff for sphere α • R_{MT}^{α} radius for muffin-tin sphere α • E_{l}^{α} energy parameter for u_{l}^{α} , \dot{u}_{l}^{α}

Note: LAPW basis depends on atom positions and features discontinuities at MT boundaries





Orthogonality of LAPW basis functions to core electron states (1)

• $u_l(r)$, $u_l^c(r)$ given by radial Schrödinger equation:

$$\left[-\frac{1}{2}\frac{\partial^2}{\partial r^2} + \frac{l(l+1)}{2r^2} + V_{\text{eff}}^{\text{sphr}}(r)\right]ru_l(r) = E_l ru_l(r)$$
(1)

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Multiply (1) for u_l(r) by ru^c_l(r) and vice versa, subtract the two resulting equations from each other, and integrate:

$$\int_{0}^{R_{\text{MT}}} -\frac{1}{2} r u_{l}^{c}(r) \frac{\partial^{2}}{\partial r^{2}} r u_{l}(r) + \frac{1}{2} r u_{l}(r) \frac{\partial^{2}}{\partial r^{2}} r u_{l}^{c}(r) dr = (E_{l} - E_{l}^{c}) \int_{0}^{R_{\text{MT}}} u_{l}(r) r^{2} u_{l}^{c}(r) dr$$

$$= \text{Assumption: } u_{l}^{c}(r)|_{R_{\text{MT}}} = 0, \ \frac{\partial}{\partial r} u_{l}^{c}(r)|_{R_{\text{MT}}} = 0$$

$$= \text{We obtain: } 0 = \left\langle u_{l}^{c} | u_{l} \right\rangle_{R_{\text{MT}}} \text{ and analogously } 0 = \left\langle u_{l}^{c} | \dot{u}_{l} \right\rangle_{R_{\text{MT}}}$$



Orthogonality of LAPW basis functions to core electron states (2)

- Orthogonality allows to determine core and valence electron energies and wave functions separately from each other
- Core electrons
 - Representation for each atom separately on radial mesh
 - Fully relativistic treatment
- Valence electrons
 - Representation by LAPW basis
 - Scalar-relativistic description in MT spheres
 - Optional inclusion of spin-orbit coupling
- But: assumption $u_l^c(r)|_{R_{MT}} = 0$, $\frac{\partial}{\partial r}u_l^c(r)|_{R_{MT}} = 0$ only approximately fulfilled
 - Semicore states can lead to ghost bands





The linearization within the LAPW basis

 Description in MT spheres is not systematically improved by increasing the reciprocal cutoff parameter K_{max}

Linearization of solutions u_l at arbitrary energy ϵ

- $= u_l^{\alpha}(r_{\alpha},\epsilon) = u_l^{\alpha}(r_{\alpha},E_l^{\alpha}) + (\epsilon E_l^{\alpha})\dot{u}_l^{\alpha}(r_{\alpha},E_l^{\alpha}) + \mathcal{O}\left[(\epsilon E_l^{\alpha})^2\right]$
- Due to the restriction to the function space spanned by $u_l^{\alpha}(r_{\alpha}, E_l^{\alpha})$ and $\dot{u}_l^{\alpha}(r_{\alpha}, E_l^{\alpha})$ we obtain a linearization error.
- This description is sufficient to obtain accurate results for many materials.





Extending the LAPW basis with local orbitals

Additional basis functions localized in MT spheres

 $\phi_{L}^{\mathsf{lo}}(\mathbf{r}) = \left[a_{L}^{\mathsf{lo}}u_{I}^{\alpha}(\mathbf{r}_{\alpha}, \mathbf{E}_{I}^{\alpha}) + b_{L}^{\mathsf{lo}}\dot{u}_{I}^{\alpha}(\mathbf{r}_{\alpha}, \mathbf{E}_{I}^{\alpha}) + c_{L}^{\mathsf{lo}}u_{I}^{\alpha}(\mathbf{r}_{\alpha}, \mathbf{E}_{I}^{\mathsf{lo}})\right]Y_{L}(\hat{\mathbf{r}}_{\alpha})$

- Mainly used to describe semicore states
- Determination of a^{lo}_L, b^{lo}_L, and c^{lo}_L by enforcing zero value and slope at the MT boundary, as well as a normalization condition on the local orbital





Extending the LAPW basis with local orbitals

Additional basis functions localized in MT spheres

 $\phi_{L}^{\mathsf{lo}}(\mathbf{r}) = \left[a_{L}^{\mathsf{lo}}u_{I}^{\alpha}(\mathbf{r}_{\alpha}, \mathbf{E}_{I}^{\alpha}) + b_{L}^{\mathsf{lo}}\dot{u}_{I}^{\alpha}(\mathbf{r}_{\alpha}, \mathbf{E}_{I}^{\alpha}) + c_{L}^{\mathsf{lo}}u_{I}^{\alpha}(\mathbf{r}_{\alpha}, \mathbf{E}_{I}^{\mathsf{lo}})\right]Y_{L}(\hat{\mathbf{r}}_{\alpha})$

- Semicore states (SCLO)
 - Choose E^{lo}_l to be energy of semicore state
- Unoccupied orbitals (HELO)
 - Choose *E*^{lo} above Fermi energy
- Higher derivative LOs (HDLO)
 - Choose ü_l^α(r_α, E_l^α) instead of u_l^α(r_α, E_l^o)





The LAPW basis for films

$$\phi_{\mathbf{k}_{\parallel}\mathbf{G}}(\mathbf{r}) = \begin{cases} \frac{1}{\sqrt{\Omega}} e^{i(\mathbf{k}_{\parallel}+\mathbf{G})\mathbf{r}} & \text{for } \mathbf{r} \in \mathbf{IR} \\ \sum_{L} \left[a_{\mathbf{k}_{\parallel}\mathbf{G}}^{L\alpha} u_{l}^{\alpha}(r_{\alpha}, E_{l}^{\alpha}) + b_{\mathbf{k}_{\parallel}\mathbf{G}}^{L\alpha} \dot{u}_{l}^{\alpha}(r_{\alpha}, E_{l}^{\alpha}) \right] Y_{L}(\hat{\mathbf{r}}_{\alpha}) & \text{for } \mathbf{r} \in \mathbf{MT}^{\alpha} \\ \left[a_{\mathbf{k}_{\parallel}\mathbf{G}}^{\text{vac}} u_{\mathbf{k}_{\parallel}\mathbf{G}_{\parallel}}^{\text{vac}}(z, E^{\text{vac}}) + b_{\mathbf{k}_{\parallel}\mathbf{G}}^{\text{vac}} \dot{u}_{\mathbf{k}_{\parallel}\mathbf{G}_{\parallel}}^{\text{vac}}(z, E^{\text{vac}}) \right] \\ \times \frac{1}{\sqrt{A}} e^{i(\mathbf{k}_{\parallel}+\mathbf{G}_{\parallel})\mathbf{r}_{\parallel}} & \text{for } \mathbf{r} \in \mathbf{VR}^{\text{vac}} \end{cases}$$



Vacuum Region 1 (VR¹)

Interstitial Region (IR)

Muffin-tin (MT)

Vacuum Region 2 (VR²)

• $u_{\mathbf{k}_{\parallel}\mathbf{G}_{\parallel}}^{\text{vac}}$, $\dot{u}_{\mathbf{k}_{\parallel}\mathbf{G}_{\parallel}}^{\text{vac}}$: solutions, energy derivatives to vacuum potential at energy parameters E^{vac}

•
$$G_{\perp} = 2\pi n/\tilde{D}$$



THE EXASCALE

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The LAPW basis for films

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Vacuum Region 1 (VR¹)

Interstitial Region (IR)

Muffin-tin (MT)

Vacuum Region 2 (VR²)

Parameters:

- D vacuum boundary
- *D* determination of *G*_⊥
- E^{vac} vacuum energy parameters





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Representation of density and potential

Plane-wave part

$$ho^{\mathsf{PW}}(\mathbf{r}) = \sum\limits_{\mathbf{G}}^{G_{\mathsf{max}}}
ho_{\mathbf{G}} \cdot \boldsymbol{e}^{i\mathbf{Gr}}$$

- Actually represented by stars
 - Linear combinations of plane waves according to symmetry

 $\begin{aligned} & \text{MT sphere } \alpha \\ \rho^{\alpha}(\mathbf{r}) = \sum_{L}^{L_{\text{max}}^{\alpha}} \rho_{L}^{\alpha}(r_{\alpha}) Y_{L}(\hat{\mathbf{r}}_{\alpha}) \end{aligned}$

- Actually represented by lattice harmonics
 - Linear combinations of spherical harmonics according to symmetry

Parameters:

- G_{max} reciprocal plane-wave cutoff for density, potential
- G_{maxXC} reduced cutoff for exchange correlation potential





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Choice of K_{max} , I^{α}_{max} , and $I^{\alpha}_{nonsphr}$

• Rayleigh expansion of planes waves at MT boundary:



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Choice of G_{max} and G_{maxXC}

- G_{max} is cutoff for different functions
 - Plane wave part of charge density \(\rho^{PW}(r)\)
 - Plane wave part of potential $V_{\text{eff}}^{\text{PW}}(\mathbf{r})$
 - Step function $\Theta(\mathbf{r})$ indicating the interstitial region
- $V_{\text{eff}}^{\text{PW}}(\mathbf{r})$ and $\Theta(\mathbf{r})$ have infinitely many plane-wave coefficients.
- Interstitial potential contribution to Hamilton matrix:

$$\left\langle \phi_{\mathbf{kG}} \middle| \Theta(\mathbf{r}) V_{\mathsf{eff}}^{\mathsf{PW}}(\mathbf{r}) \middle| \phi_{\mathbf{kG}'} \right
angle$$

Rule

 $G_{\max} \geq G_{\max XC} \geq 2 \cdot K_{\max}$

• typically $G_{max} \approx 3 \cdot K_{max}$, $G_{maxXC} \approx 2.5 \cdot K_{max}$





Choice of the energy parameters

fcc Ce

- energy center of mass of the *I*-projected DOS
 - minimizes quadratic error weighted by charge in each eigenstate
- atomic solutions
 - yields more friendly convergence behavior





Choice of MT radii

- Due to different bonding lengths in different materials the R_{MT} are material dependent.
- If calculations have to be compared identical MT radii should be chosen.

Large MT radii

- Faster calculations
- Larger linearization error
- Fewer SCLOs needed
- Some quantities only evaluated in MT

Small MT radii

- Slower calculations
- More stable calculations
- Smaller linearization error
- More SCLOs needed
- More space available for structural relaxations







■ *R*_{MT} = 2.25 *a*₀, lostElectrons = 0.086





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■ *R*_{MT} = 2.20 *a*₀, lostElectrons = 0.100





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■ *R*_{MT} = 2.17 *a*₀, lostElectrons = 0.109







■ *R*_{MT} = 2.16 *a*₀, lostElectrons = 0.112





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- *R*_{MT} = 2.15 *a*₀, lostElectrons = 0.123
- This error message can also have other causes.
- Other error messages are also possible.





Semicore states and ghost bands - with SCLO



• $R_{\text{MT}} = 2.16 a_0$, lostElectrons = 0.012, SCLO for 3*p* state

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Linearization error depending on energy mismatch

fcc Cerium

bcc Vanadium



• $\Delta_l = \sqrt{\parallel u_l(r,\epsilon) - \tilde{u}_l(r,\epsilon) \parallel^2}$



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The linearization error and MT radii fcc Ce rock-salt KCI



 lattice constant changes by 1.6% when MT radius is reduced lattice constant changes by 0.4% when MT radii are reduced





The linearization error and unoccupied states





 KS band gap for rock-salt KCl is reduced by 4% by adding one set of HDLOs



 KS band gap for fcc Ar is reduced by 19% by adding one set of HDLOs

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The input file generator

- Fleur uses complex input
 - Input file with default parameters is generated by input file generator inpgen
 - Simple structural input needed for inpgen
- Example input for inpgen

```
NaCl bulk

&lattice latsys='fcc' a0=10.62026 /

2

11 0.0 0.0 0.0

17 0.5 0.5 0.5
```

Usage: inpgen < myInputFile.txt</p>





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Input file generator - command line options

- -h write out list of command line options
- Optional Input to be generated
 - -explicit add some optional input directly to inp.xml file.
 - List of k points
 - Symmetry operations (otherwise if needed in sym.out file)
 - Noncollinear magnetism input
 - -noco generate noncollinear magnetism input in inp.xml
 - -genEnpara generate enpara file for more energy parameter options
 - ...





Fleur input: the inp.xml file

- Usage: fleur
- -h command line option to display all fleur modes





The inp.xml file - calculationSetup

```
<calculationSetup>
<cutoffs Kmax="3.6" Gmax="10.7" GmaxXC="8.9"
numbands="0"/>
<scfLoop itmax="15" minDistance=".0"
imix="Anderson" alpha=".05"/>
<coreElectrons ctail="T" frcor="F" kcrel="0"/>
<magnetism jspins="1" l_noco="F"/>
<bzIntegration valenceElectrons="16.0" mode="hist"
fermiSmearingEnergy=".001">
<kPointCount count="15" gamma="F"/>
</bzIntegration>
</calculationSetup>
```

+ Optional input, e.g., nocoParams, geometryOptimization





The inp.xml file - cell, xcFunctional

```
<cell>
<symmetryFile filename="sym.out"/>
<bulkLattice scale="1.000" latnam="any">
<bravaisMatrix>
<row-1> 0.00000 5.31013 5.31013</row-1>
<row-2> 5.31013 0.00000 5.31013</row-2>
<row-3> 5.31013 5.31013 0.00000</row-3>
</bravaisMatrix>
</bulkLattice>
</cell>
<xcFunctional name="pbe"
relativisticCorrections="F"/>
```





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The inp.xml file - atomSpecies

```
<atomSpecies>
   <species name="Na-1" element="Na" atomicNumber="11"</pre>
            coreStates="1" magMom=".00" flipSpin="T">
      <mtSphere radius="2.80" gridPoints="925"
                logIncrement=".0120"/>
      <atomicCutoffs lmax="10" lnonsphr="8"/>
      <energyParameters s="3" p="3" d="3" f="4"/>
      <lo type="SCLO" 1="0" n="2" eDeriv="0"/>
      <lo type="SCLO" l="1" n="2" eDeriv="0"/>
   </species>
   <species name="Cl-2" element="Cl" atomicNumber="17"</pre>
            coreStates="4" magMom=".00" flipSpin="T">
   </species>
</atomSpecies>
```



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The inp.xml file - atomGroups

+ Optional input, e.g., nocoParams





The inp.xml file - output

+ Optional input, e.g., vacuumDOS, plotting, chargeDensitySlicing

Full documentation of inp.xml file on www.flapw.de





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Overview on files

with HDF5

- inp.xml
- (sym.out)
- enpara (optional)
- cdn.hdf
- mixing_history.*
- input for special calculations
- out, inf, out.xml

without HDF5

- inp.xml
- (sym.out)
- enpara (optional)
- cdn1, cdn??, cdnc
- mixing_history.*
- stars, wkf2
- input for special calculations
- out, inf, out.xml





Further reading

- The FLAPW method
 - Overview book Singh et al., Planewaves, Pseudopotentials, and the LAPW Method, Springer
 - Initial publication Andersen, PRB 12, 3060 (1975)
 - FLAPW for films Krakauer *et al.*, PRB **19**, 1706 (1979)
 - Potential calculation Weinert, J.Math.Phys. 22, 2433 (1981)
 - Predecessor (APW) Slater, Phys.Rev. 51, 846 (1937)
- Local orbitals
 - SCLOs Singh, PRB 43, 6388 (1991)
 - HELOs Betzinger et al., PRB 83, 045105 (2011)
 - HDLOs Friedrich et al., PRB 74, 045104 (2006)
 - Linearization error Michalicek et al., CPC 184, 2670 (2013)





Conclusions

- LAPW basis + local orbitals
- Guidelines for setting parameters
 - $K_{\text{max}}, I_{\text{max}}^{\alpha}, I_{\text{nonsphr}}^{\alpha}$
 - $R_{\rm MT}^{\alpha}, E_l^{\alpha}$
 - G_{max}, G_{maxXC}
- Semicore states and ghost bands
- The linearization error
- Fleur input files
- Not discussed
 - General numerical DFT parameters, e.g., k point set, Fermi smearing, ...







